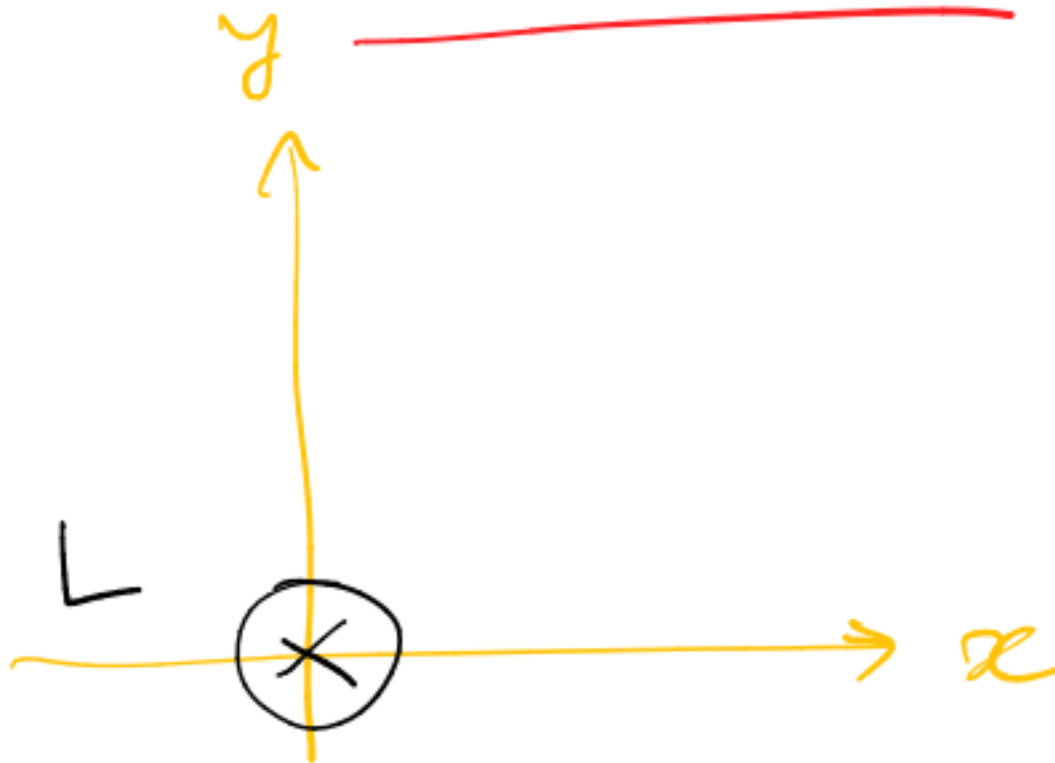


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Lecture 7



Consider
a light
wave

whose electric field
amplitudes along the

$$\begin{pmatrix} A_x e^{i\phi_x} \\ A_y e^{i\phi_y} \end{pmatrix}$$

x- & y- dims are

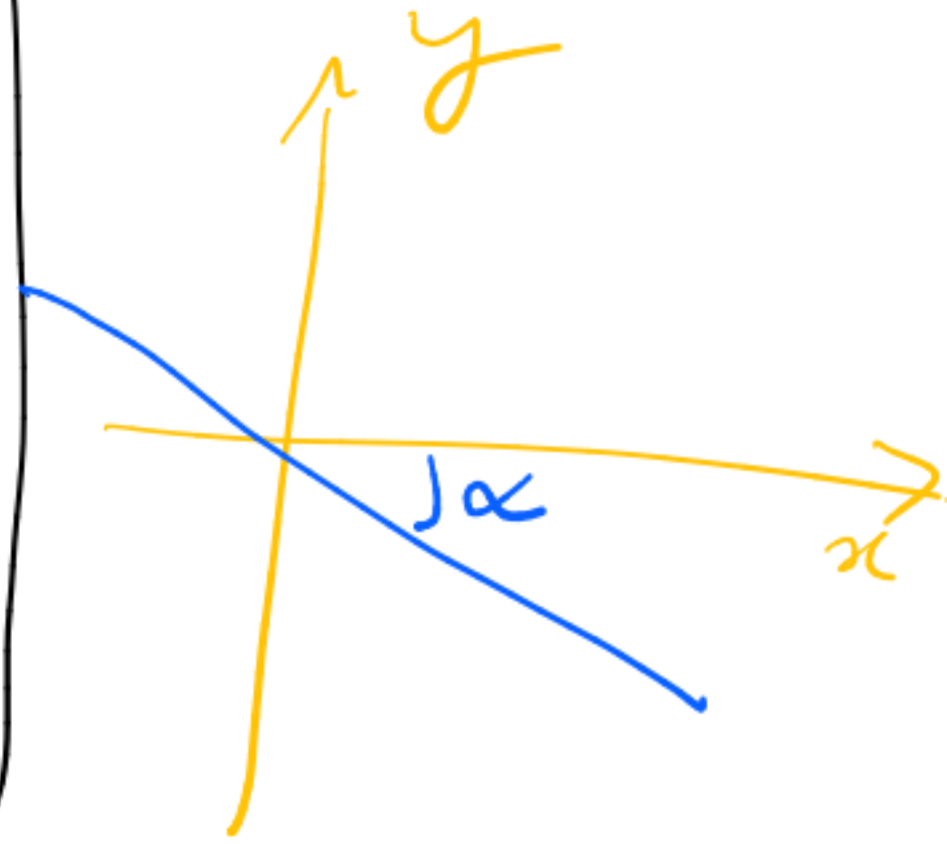
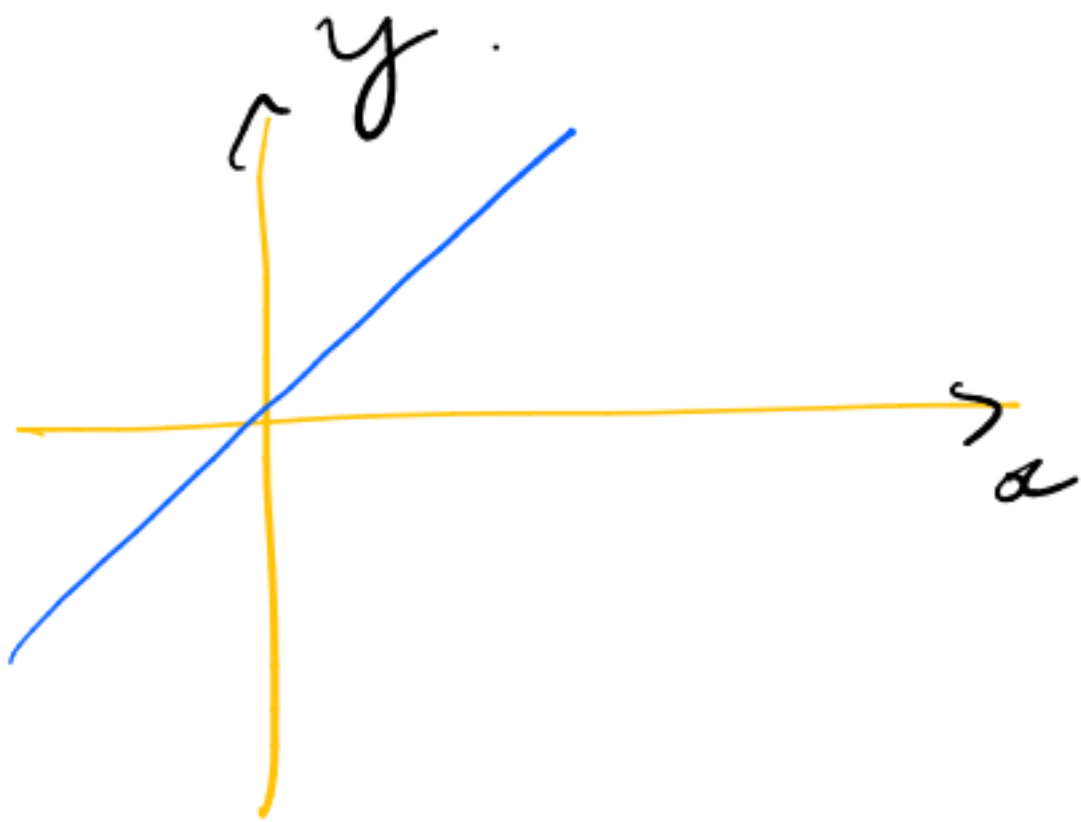
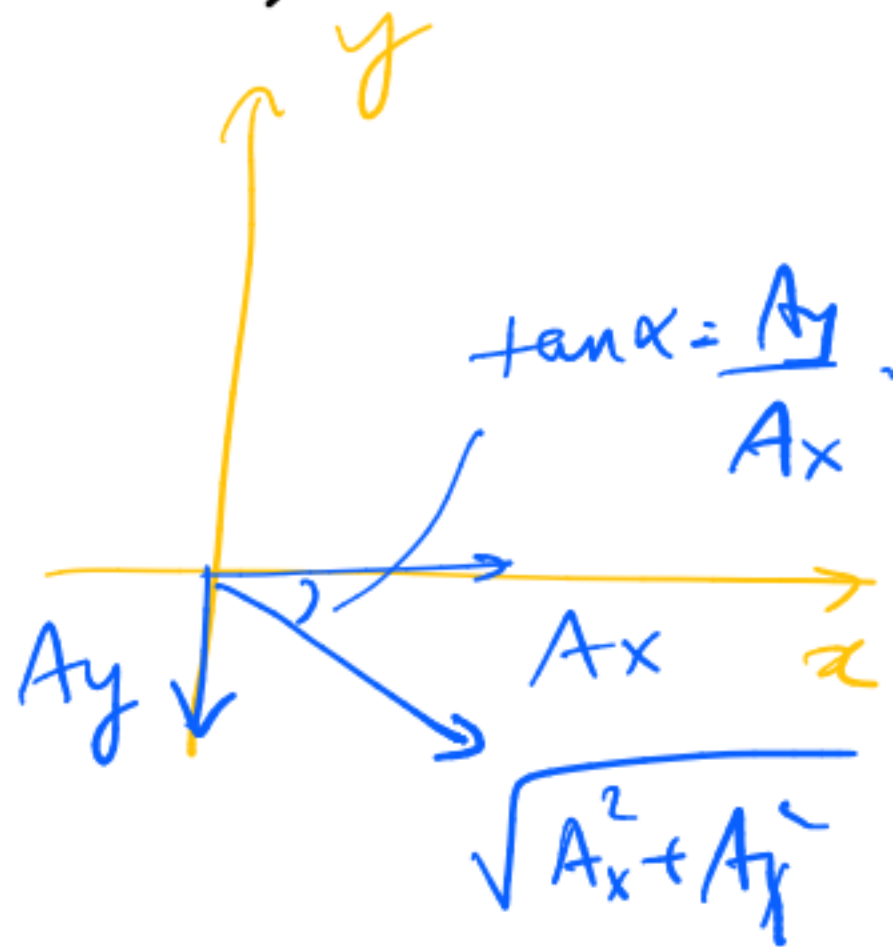
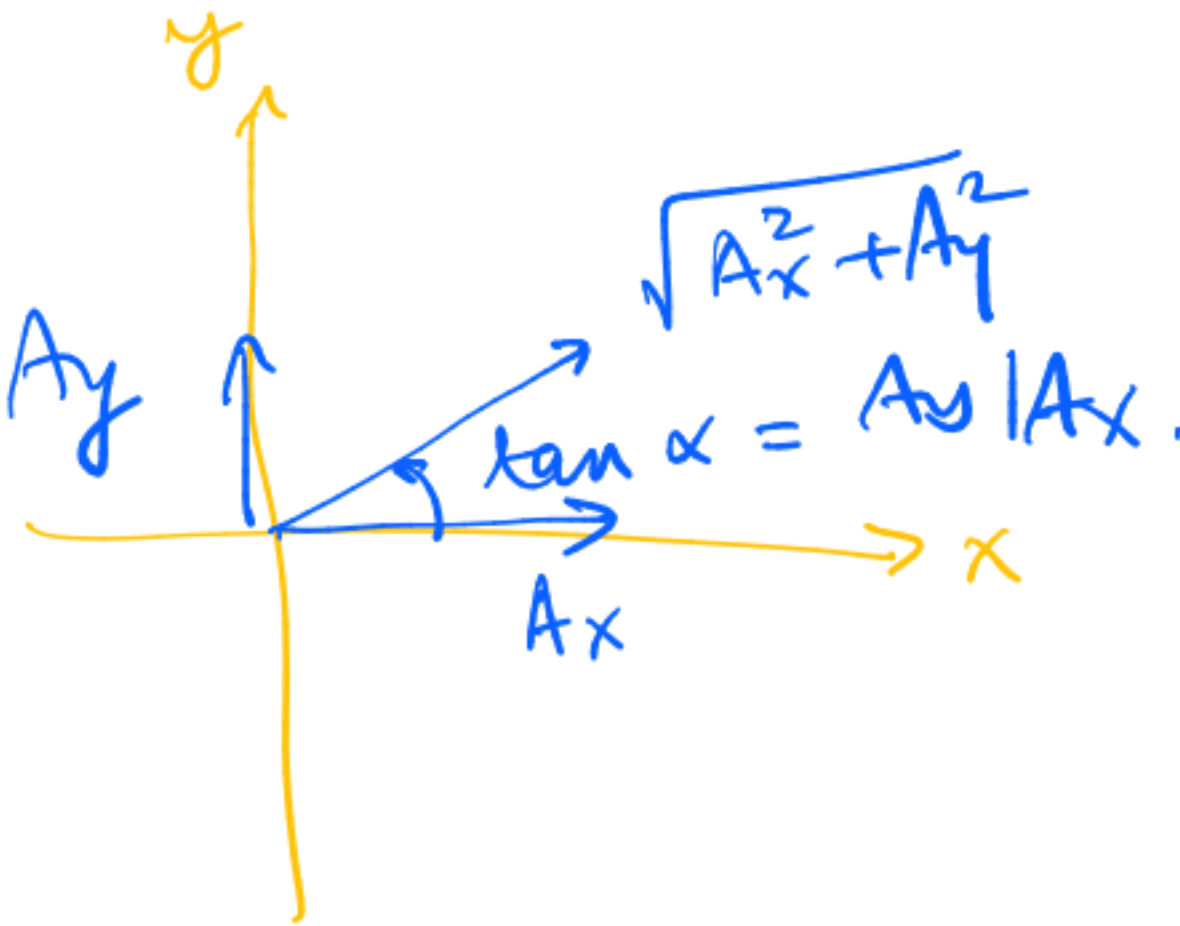
$$= \begin{pmatrix} A_x \\ A_y e^{+i\Delta} \end{pmatrix} e^{i\beta/2} e^{-i\Delta/2}$$

Plane polarised light:

$$\Delta = 0 \text{ or } \pi$$

$$\begin{pmatrix} A_x \\ A_y \end{pmatrix}$$

$$\begin{pmatrix} A_x \\ -A_y \end{pmatrix}$$





AmazonBasics Circular Polarizer Filter- 55 mm

Brand: AmazonBasics

★★★★☆ 1,311 ratings

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Circularly polarized light,

$$A_x = A_y$$

$$\Delta = \pm \pi/2$$

Light vector corresponding to circularly pol. light:

$$\begin{pmatrix} A_x \\ A_x e^{\pm i\frac{\pi}{2}} \end{pmatrix} = A_x \begin{pmatrix} 1 \\ \pm i \end{pmatrix}$$

$$E_x = \operatorname{Re} \left[A_x e^{i \left(\frac{2\pi}{\lambda} (r - ct) \right)} \right]$$

$$E_x = A_x \cos \left(\frac{2\pi}{\lambda} (r - ct) \right)$$

$$E_y = \operatorname{Re} \left[A_y e^{i \left(\frac{2\pi}{\lambda} (r - ct) \right)} \right]$$

\uparrow
 $\pm i A_x$

$$= \operatorname{Re} \left[A_x (\pm i) \left(\cos \left(\frac{2\pi}{\lambda} (r - ct) \right) + i \sin \left(\frac{2\pi}{\lambda} (r - ct) \right) \right) \right]$$

$$E_y = \mp A_x \sin \left(\frac{2\pi}{\lambda} (r - ct) \right)$$

The E_x & E_y above satisfy

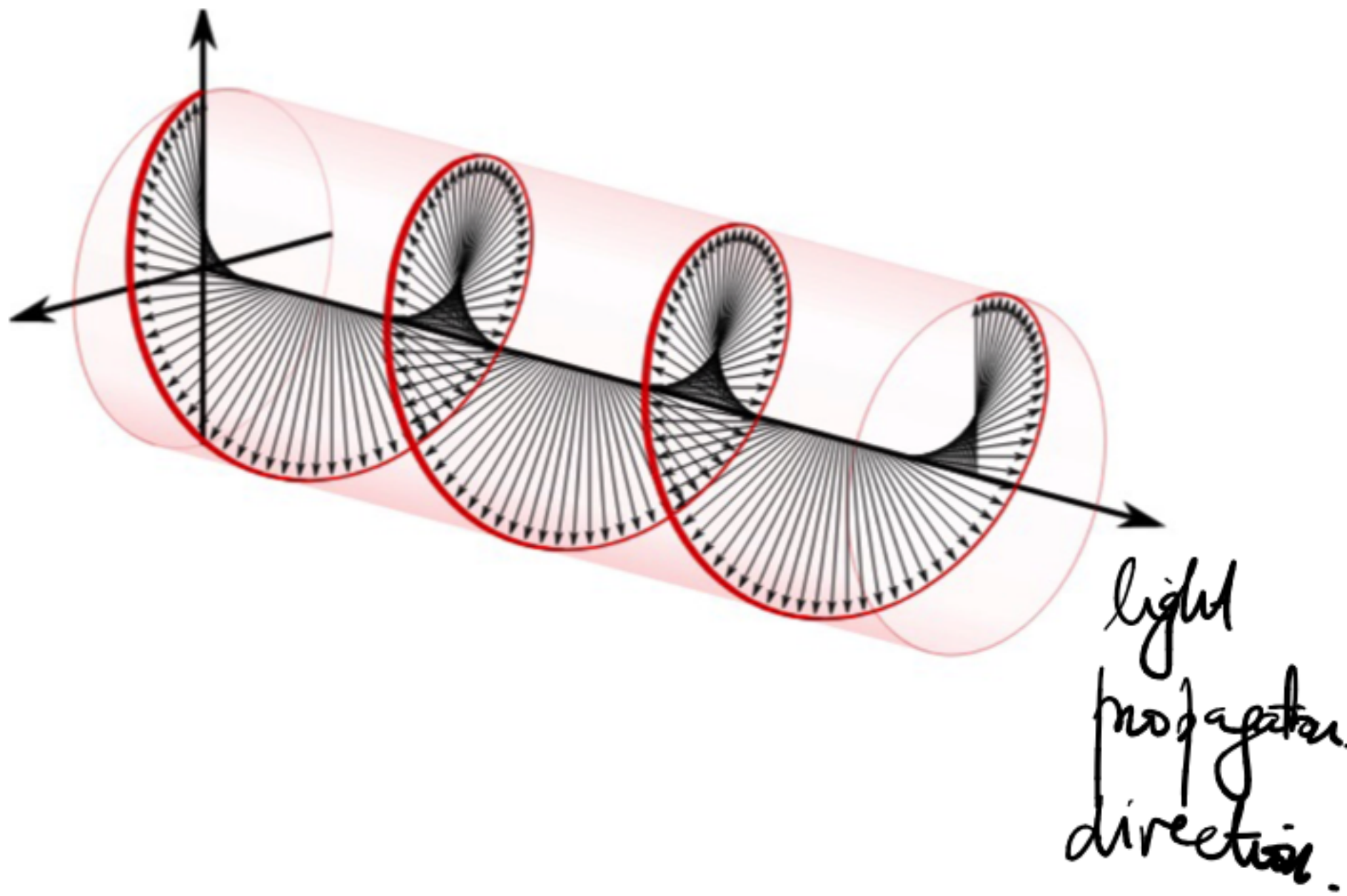
$$E_x^2 + E_y^2 = A_x^2$$

Angle that the resultant vector makes with the x -axis

$$\tan \alpha = \frac{E_y}{E_x} = \mp \tan \left(\frac{2\pi}{\lambda} (x - ct) \right)$$

$$\Rightarrow \alpha = \mp \frac{2\pi}{\lambda} (x - ct)$$

The angle the resultant \vec{E} of constant magnitude A_x makes w/ the x -axis changes linearly with time.



Elliptically polarised light.

Light vector amplitude components
along x - y - axes.

$$\begin{pmatrix} A_x \\ \pm i A_y \end{pmatrix}$$

$$E_x = A_x \cos \left(\frac{2\pi}{\lambda} (r - ct) \right)$$

$$E_y = A_y \sin \left(\frac{2\pi}{\lambda} (r - ct) \right)$$

$$\Rightarrow \frac{E_x^2}{A_x^2} + \frac{E_y^2}{A_y^2} = 1.$$

Phase angle depends also on E_y/E_x .

Plane polarizer.



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$P =$ polarising axis of the polarising sheet

Incident light wave



$\begin{pmatrix} A_x \\ A_y \end{pmatrix}$

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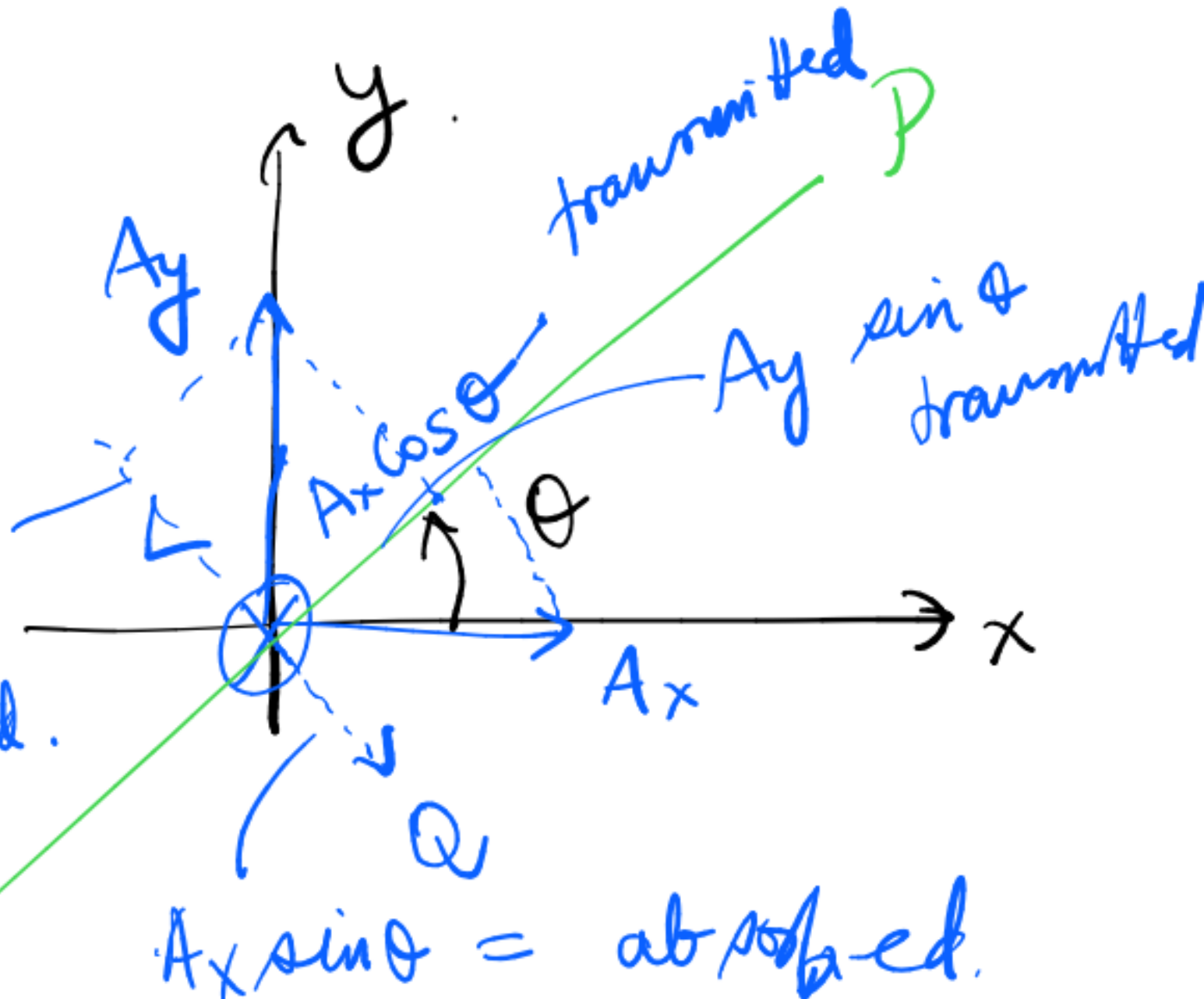
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$A_y \cos \theta$
 \downarrow
 absorbed.



Emerging light will have electric field amplitude

$$A_p = A_x \cos \theta + A_y \sin \theta$$

along the polarizer axis - P

Components of the emerging light vectors along the x- & y-

$$\text{axes} = \begin{pmatrix} A'_p \cos \theta \\ A'_p \sin \theta \end{pmatrix}$$

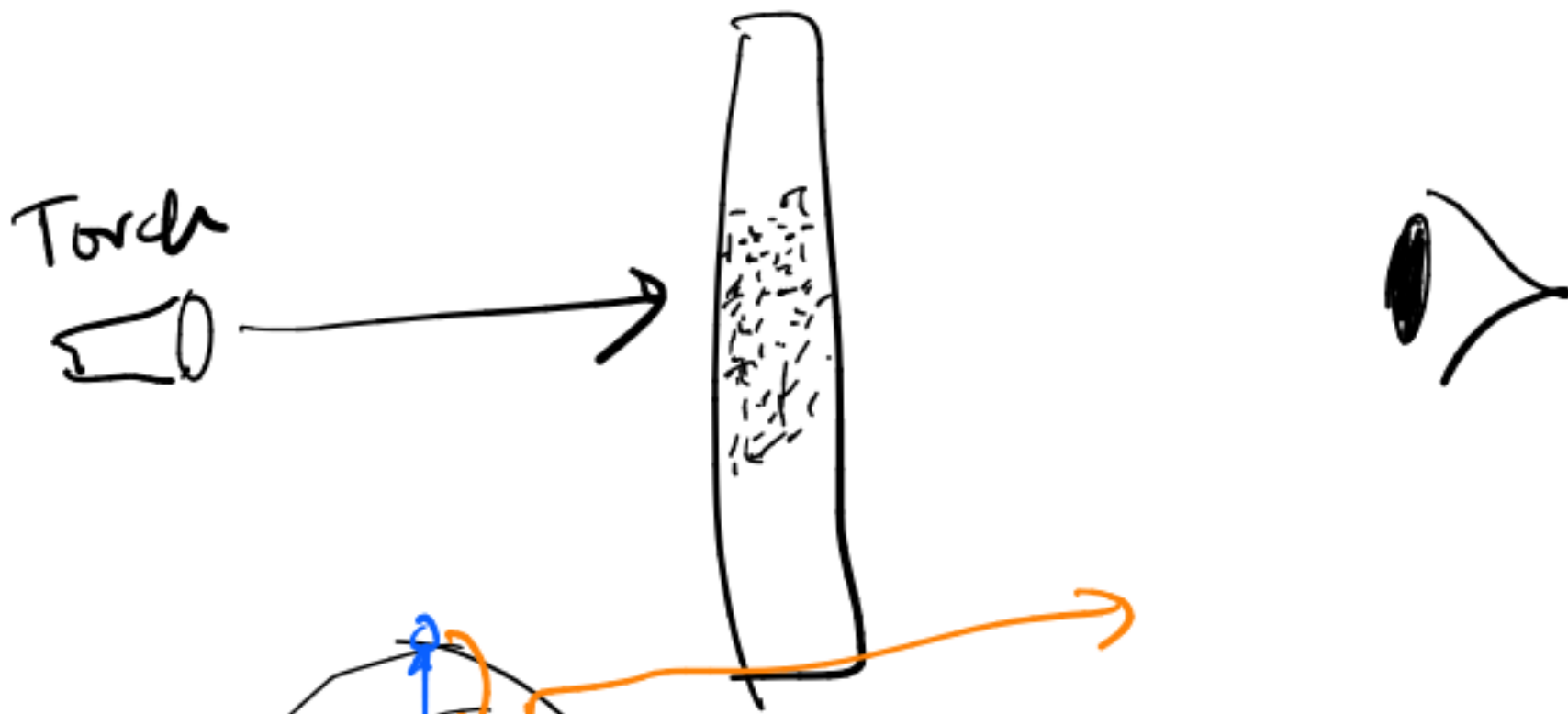
$$\Rightarrow \begin{pmatrix} A'_x \\ A'_y \end{pmatrix} = \begin{pmatrix} \cos^2 \theta & \sin \theta \cos \theta \\ \sin \theta \cos \theta & \sin^2 \theta \end{pmatrix} \begin{pmatrix} A_x \\ A_y \end{pmatrix}$$

$P_\theta =$ polarizer matrix.

Common polarizer orientation is w/
 $\theta = \pi/4$.

$$P_{\theta} = \frac{1}{2} \begin{pmatrix} 1 & & & \\ & 1 & & \\ & & 1 & \\ & & & 1 \end{pmatrix}$$

Light in media other than
vacuum.



Light interacts
with the media
it passes through.

Light travels slowly in non-vacuum media.

Speed of light in a medium = $\frac{c}{n}$,

where n = refractive index of

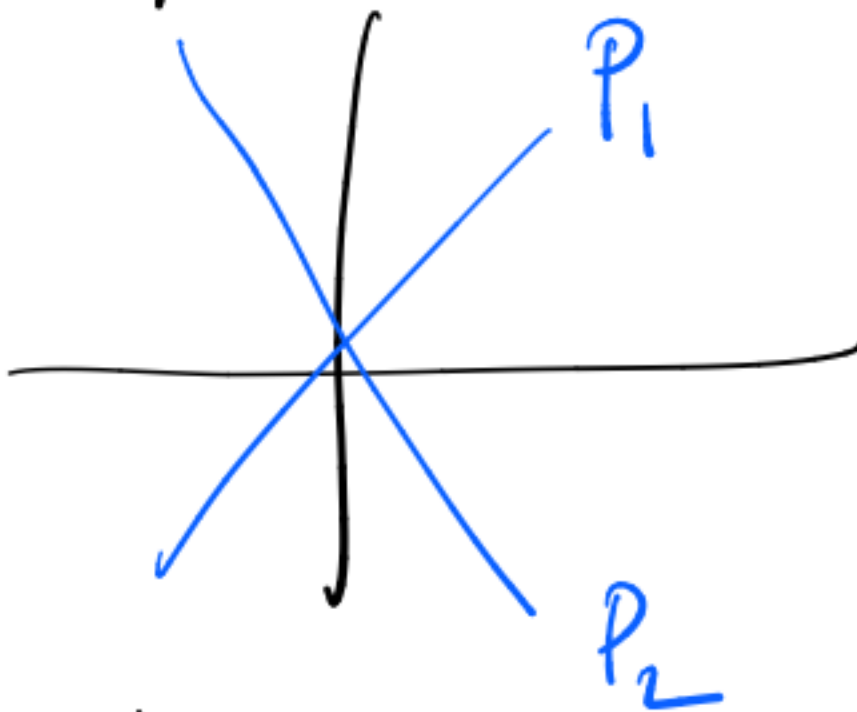
the medium.

$n = 1.3$ for water

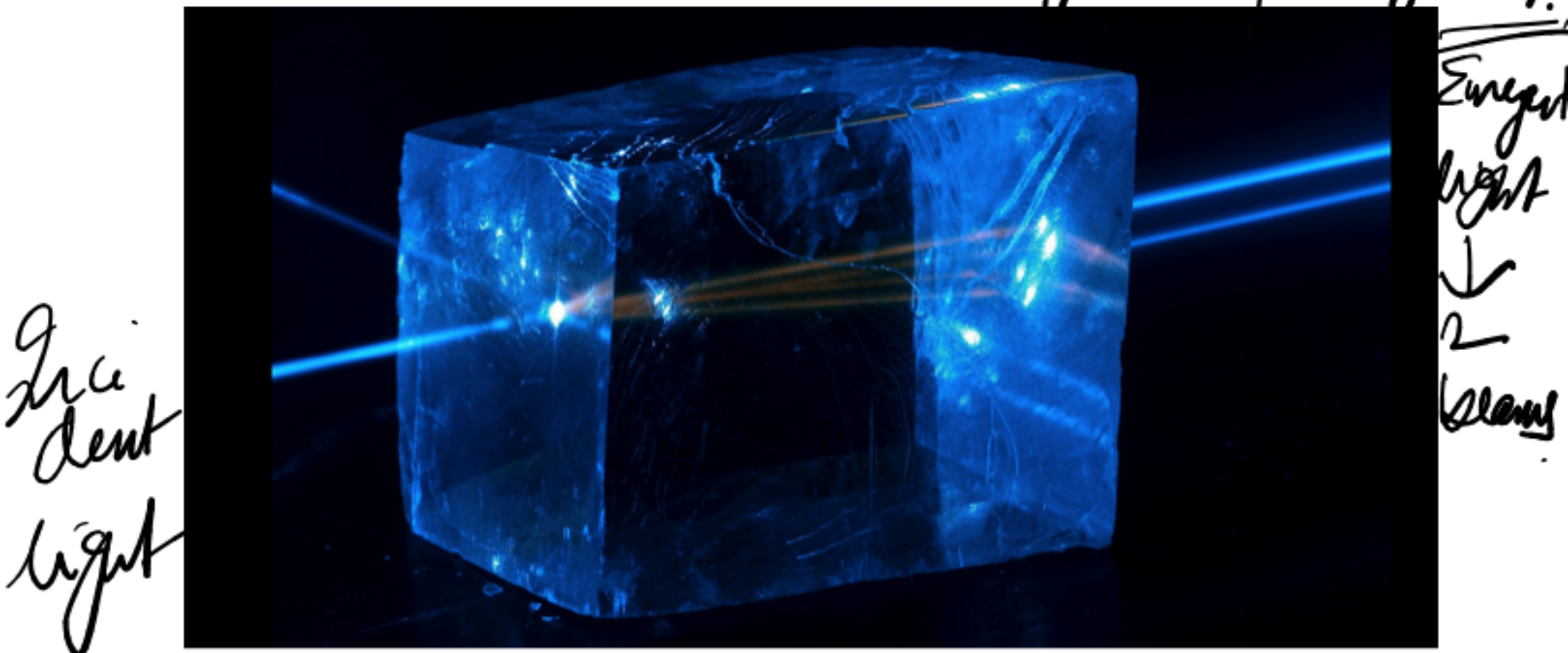
$n = 1.5$ for glass.

Birefringence.

Water is an isotropic transmitter of light \Rightarrow The speed of light in water does not depend on the polarization of light.



Calcite crystal: Speed of light/refractive index depends also on light's polarization.



The incident beam is split into multiple (2) emergent beams — Light travels through Calcite with diff. speeds depending on polarization.

Many non-crystalline transparent materials are optically isotropic (like water) when unloaded.

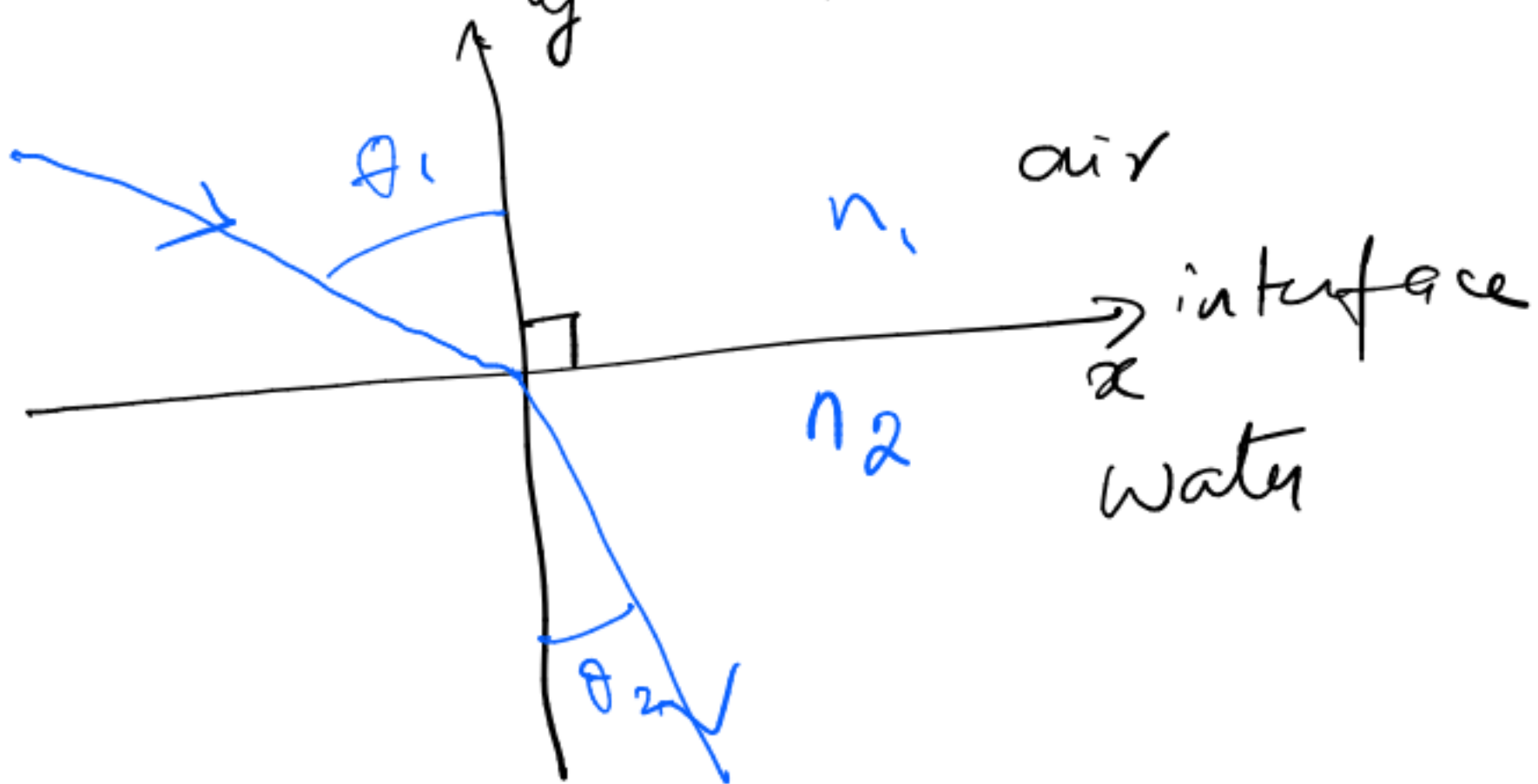
When loaded, however, they become anisotropic along the principal axes $\sigma_1, \sigma_2, \sigma_3$ of the stress tensor.

This phenomenon is called **temporary or artificial birefringence**.

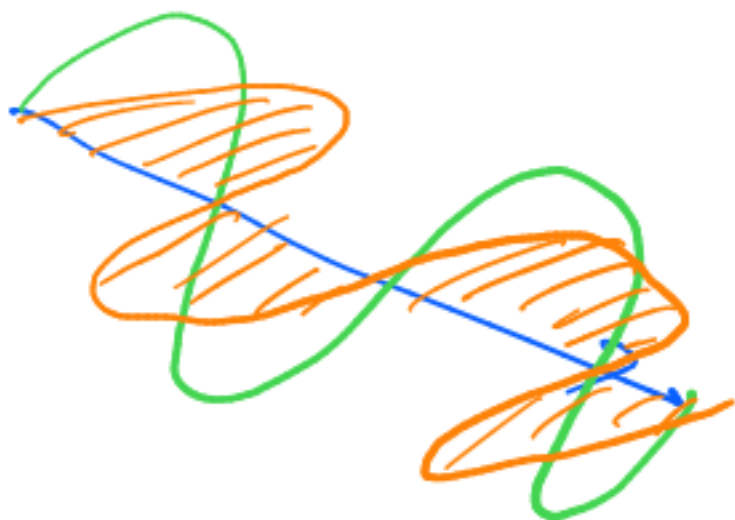
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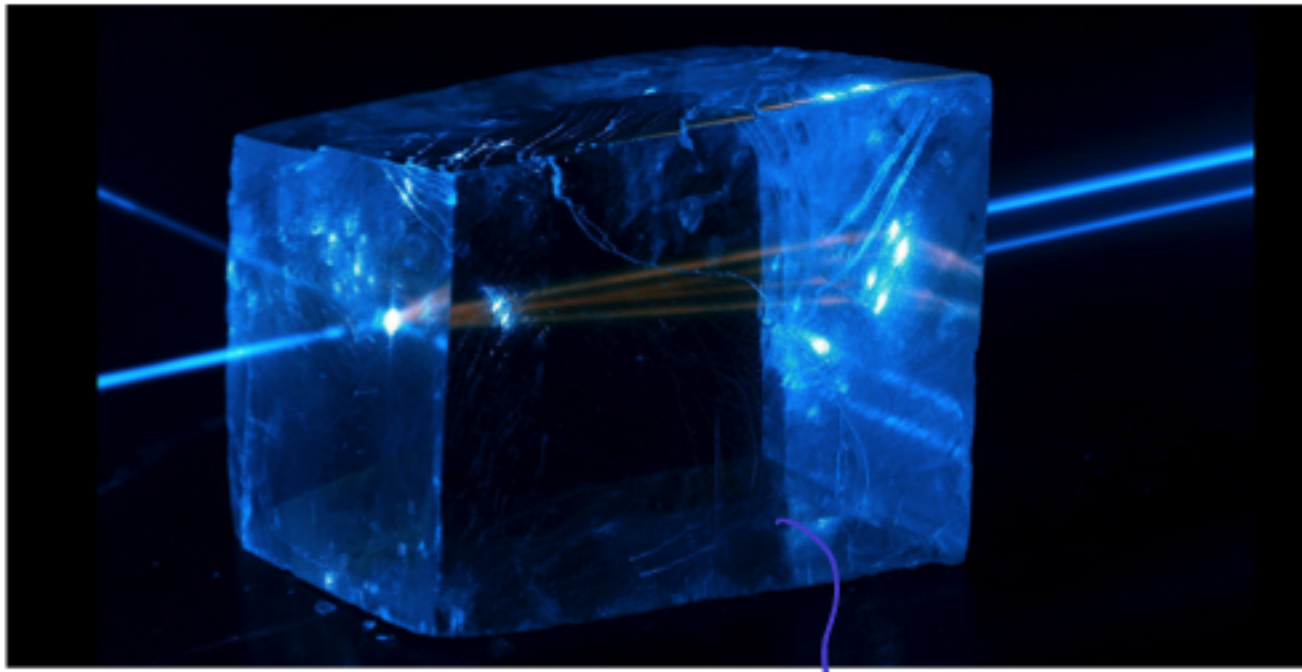
Lecture 8

Snell's law of refraction,



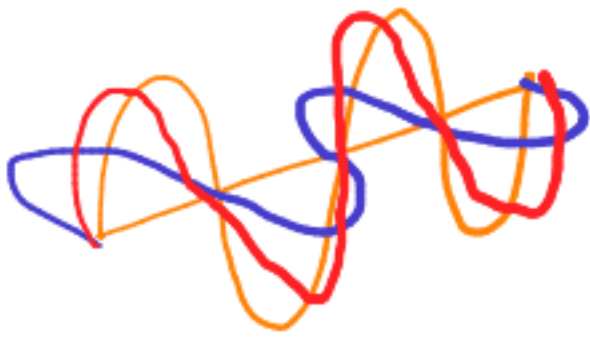
$$n_1 \sin \theta_1 = n_2 \sin \theta_2$$





Incident
laser
light
has many polarities

optically
anisotropic.



The calcite crystal is optically
anisotropic.

This comes from the anisotropy of
its crystal lattice.

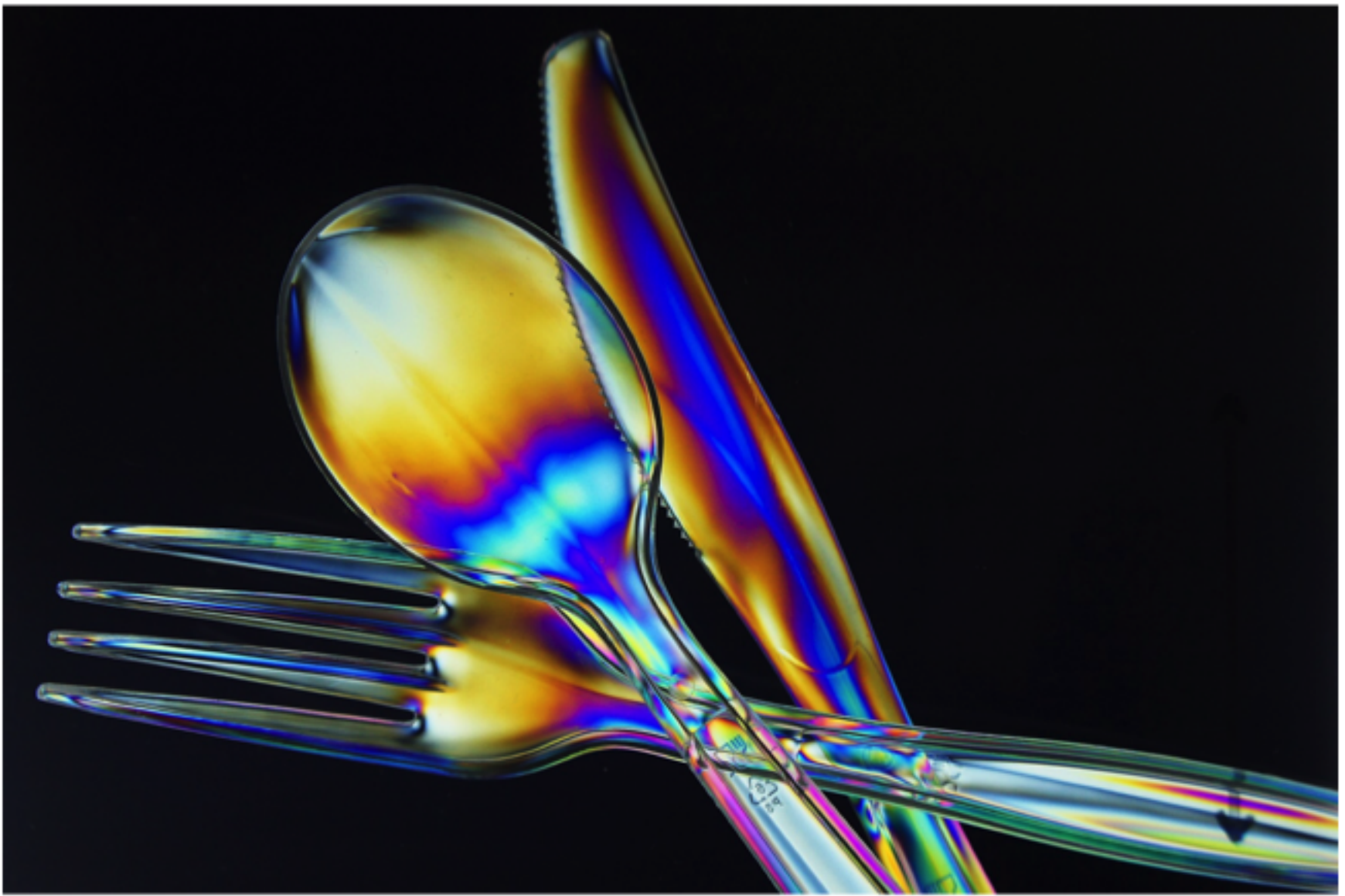
The calcite crystal displays a
different refractive index for
different polarities of the incident

light.

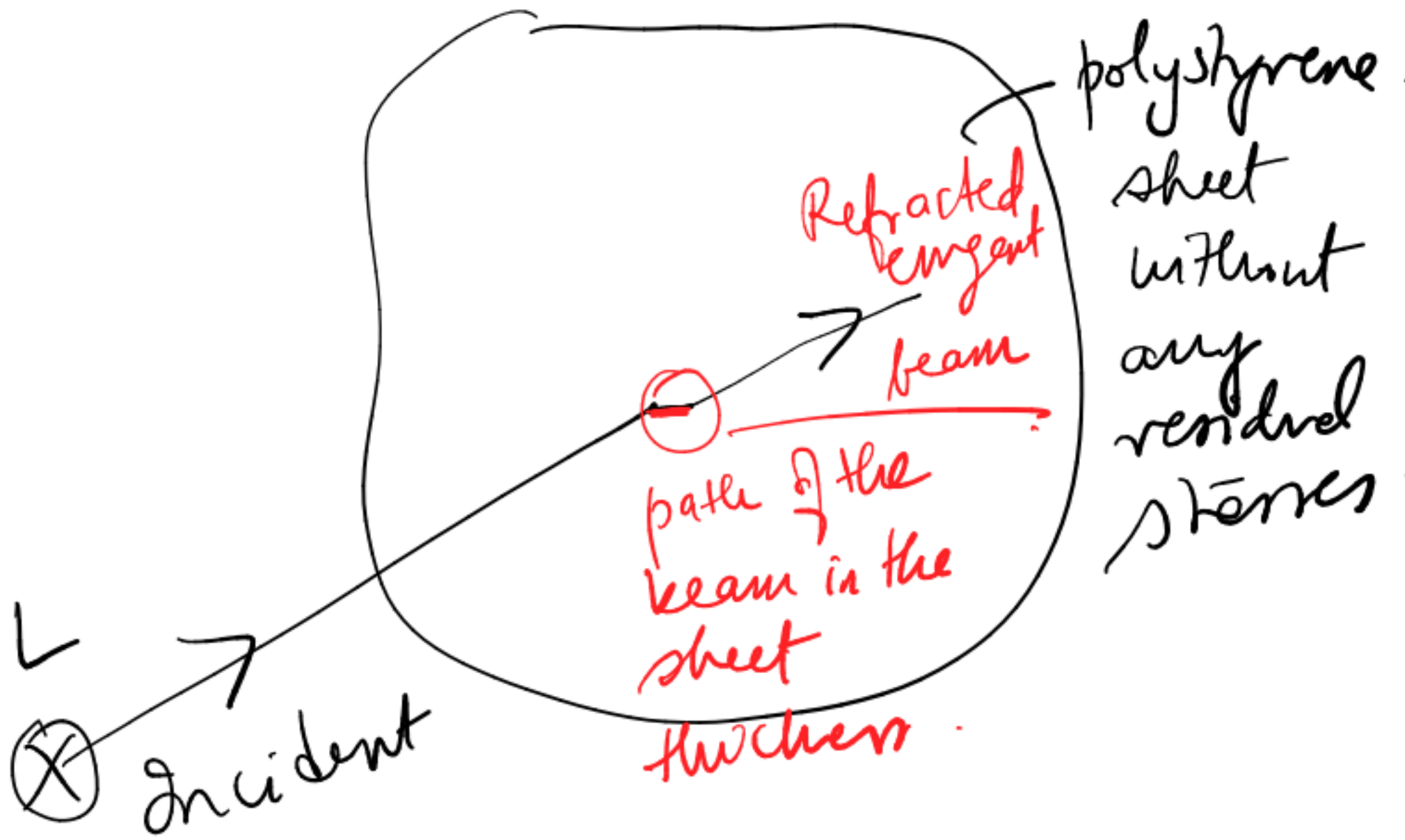
The 2 beams travel along diff paths (according to Snell's law) and at different speeds.

The property of calcite to display 2 refractive indices depending on the polarity of the incoming light beam is called double refraction or

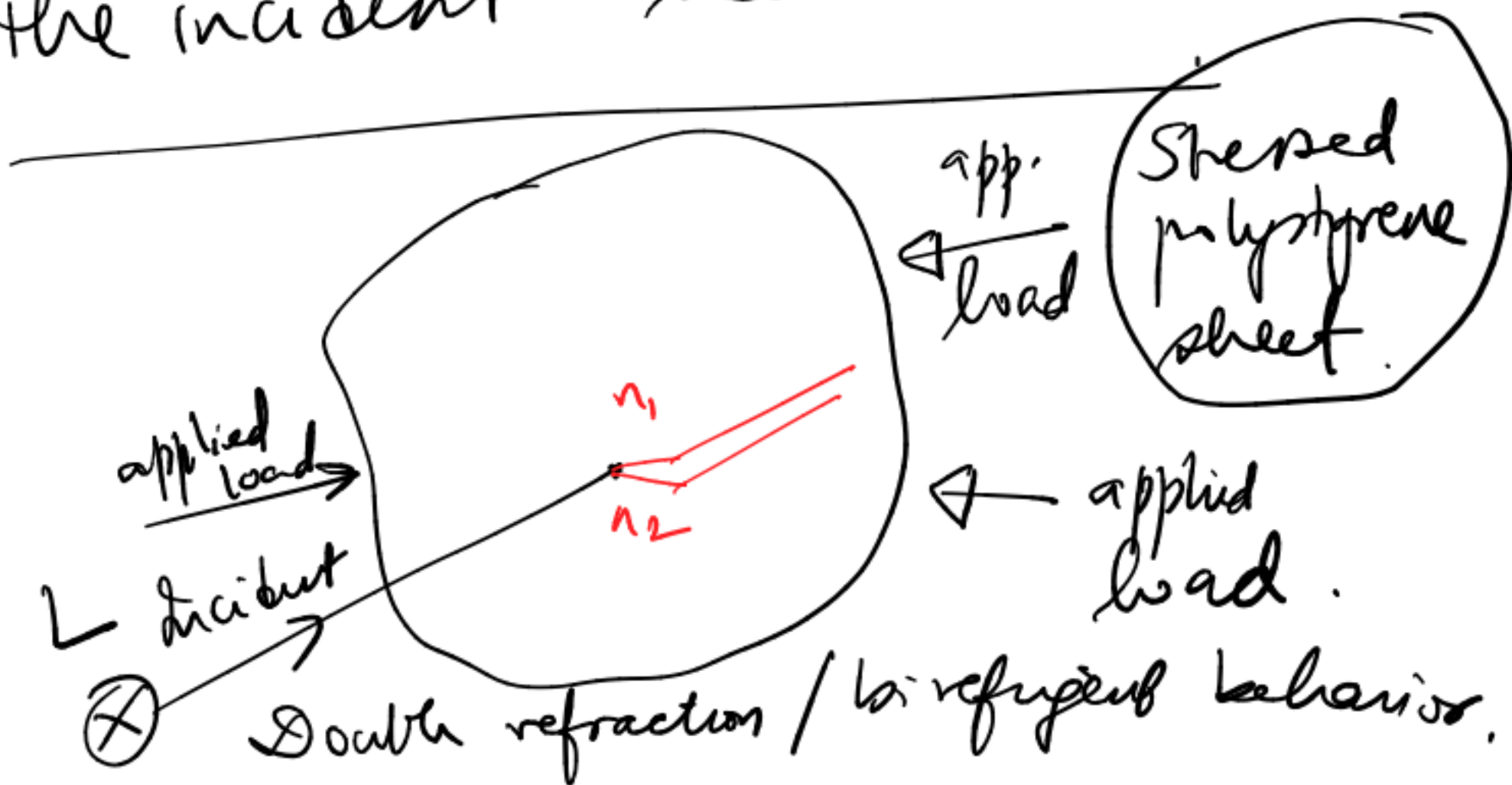
birefringence.

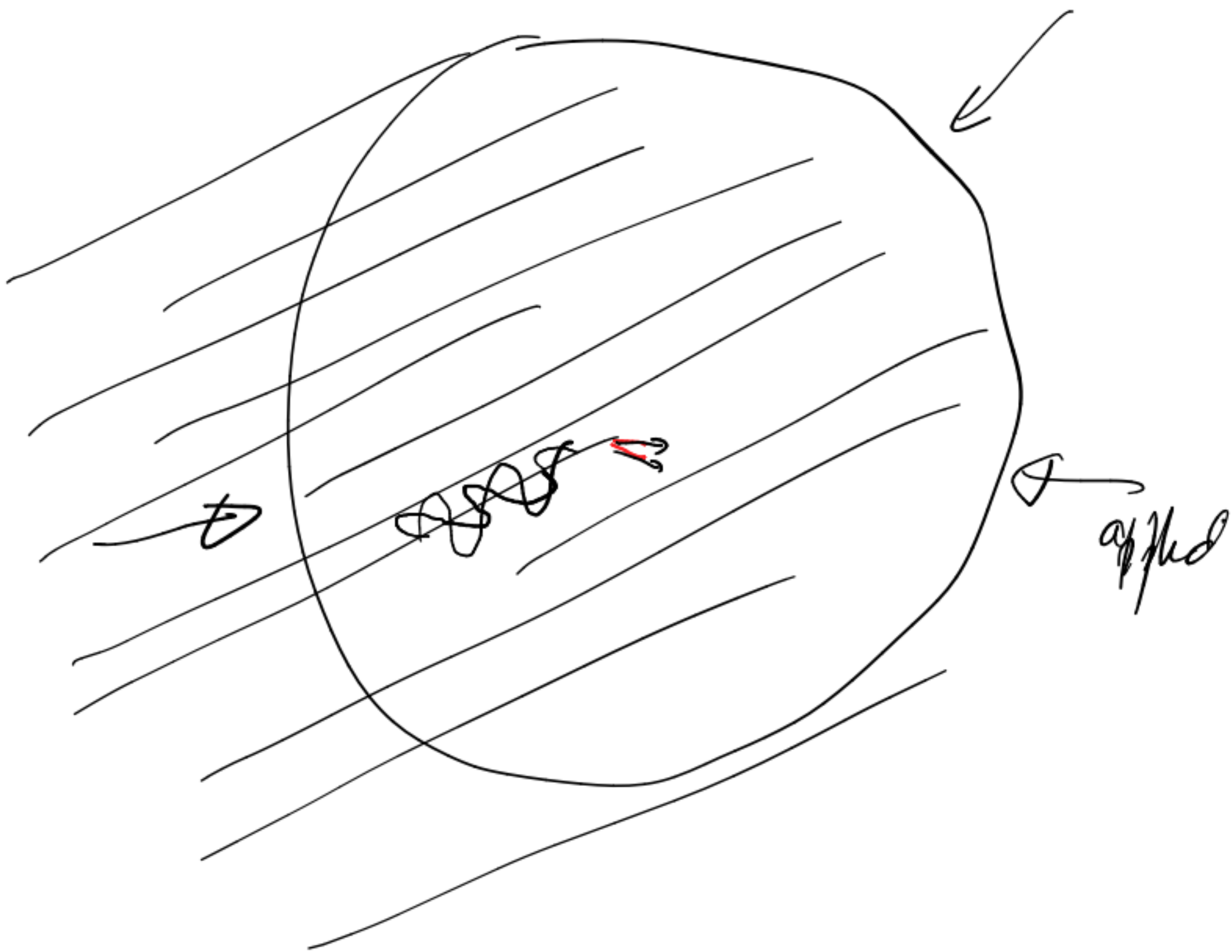


Plastics are also birefringent because of their optical anisotropy which arises during their manufacture.



Single refraction — No splitting of the incident beam.





If the entire face of the loaded plastic solid is illuminated with white / monochromatic light, we will observe fringes (color / monochromatic) on the emergent

side of the plastic birefringent plate.

Mathematically, a birefringent material introduces a phase shift, Δ between the 2 refracted waves.

Light beam passing through a birefringent sheet.

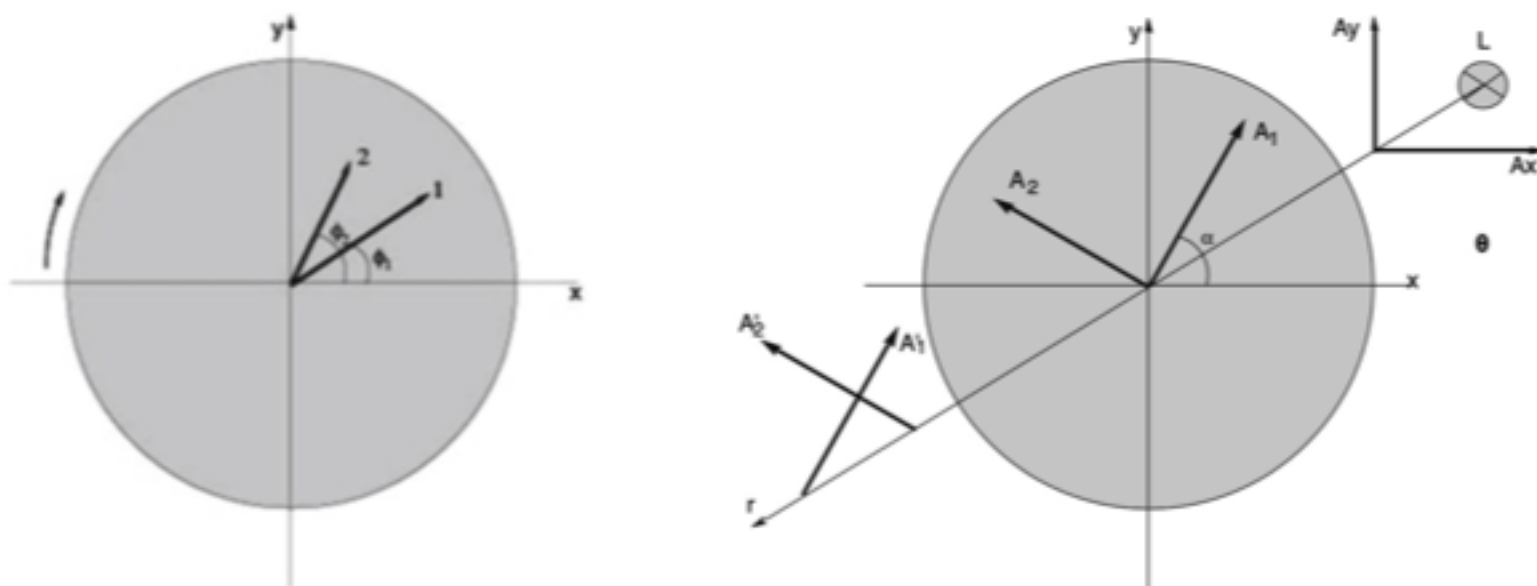


Fig. 3.6 Phase shift introduced by a birefringent filter in complex representation on the left and in physical space on the right ($\Delta = \phi_2 - \phi_1$)

Q: Given $\begin{pmatrix} A_x \\ A_y \end{pmatrix}$ are the components of the incident light beam, find $\begin{pmatrix} A'_x \\ A'_y \end{pmatrix}$ emerging from a birefringent sheet.

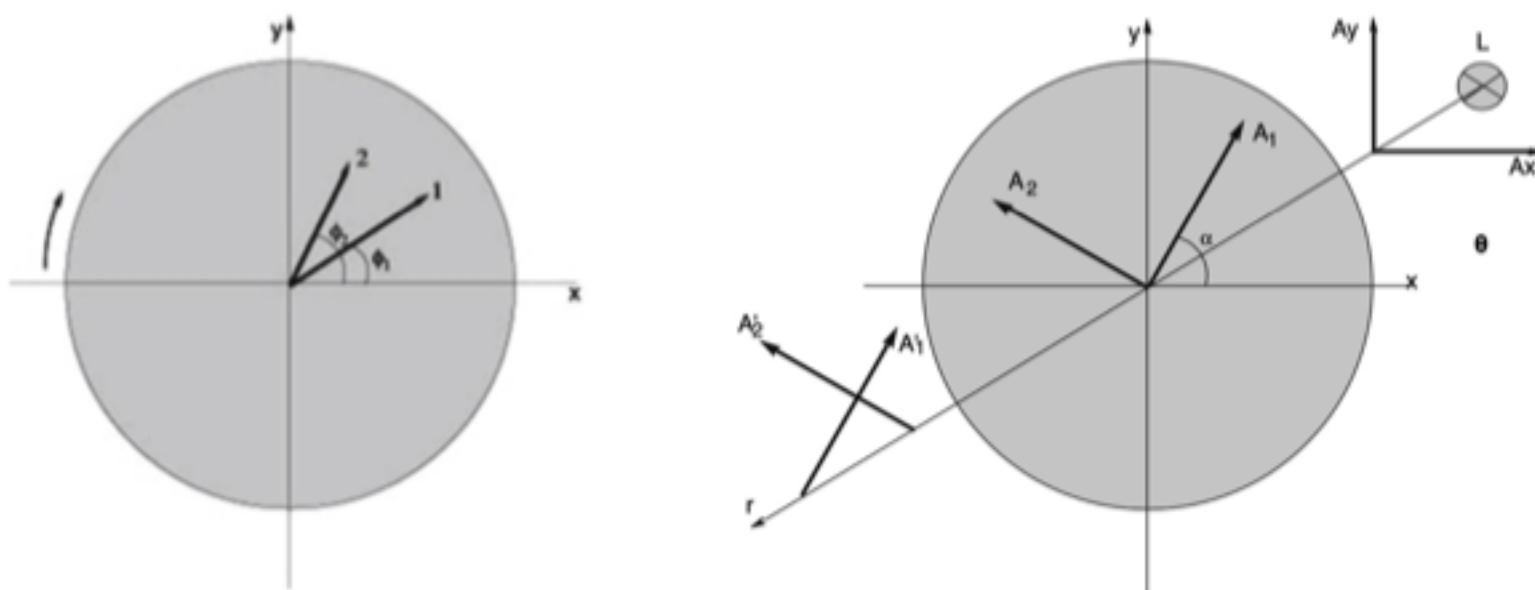
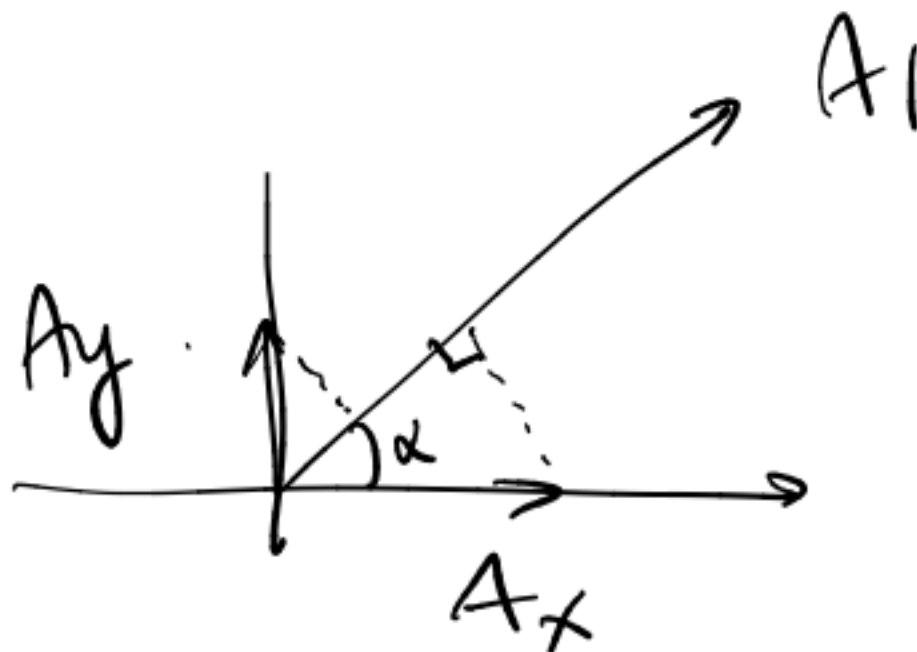


Fig. 3.6 Phase shift introduced by a birefringent filter in complex representation on the left and in physical space on the right ($\Delta = \Phi_2 - \Phi_1$)



$$A_1 = A_x \cos \alpha + A_y \sin \alpha$$

$$A_2 = -A_x \sin \alpha + A_y \cos \alpha$$

A_1, A_2 are ^{complex} components along 1, 2 axes of the incident beam.

Unlike a polarizing sheet, a birefringent sheet introduces a phase shift between the A_1 &

A_2 waves.

$$A_1' = A_1$$

$$A_2' = A_2 e^{i\Delta}$$

phase shift

introduced by diff. lights speeds for polarity along

A_1 & A_2 .

①

$$\begin{pmatrix} A_1' \\ A_2' \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & e^{i\Delta} \end{pmatrix} \begin{pmatrix} \cos\alpha & \sin\alpha \\ -\sin\alpha & \cos\alpha \end{pmatrix} \begin{pmatrix} A_x \\ A_y \end{pmatrix}$$

(2)

(1) & (2) are 2 forms of the same equation.

Rotate A_1', A_2' back to x, y coord.

system:

$$\begin{pmatrix} A_x' \\ A_y' \end{pmatrix} = \begin{pmatrix} \cos\alpha & -\sin\alpha \\ \sin\alpha & \cos\alpha \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & e^{i\Delta} \end{pmatrix} \begin{pmatrix} \cos\alpha & \sin\alpha \\ -\sin\alpha & \cos\alpha \end{pmatrix} \begin{pmatrix} A_x \\ A_y \end{pmatrix}$$

↑
transmitted
beam

↑
incident
beam

More compactly:

$$A' = M_{2 \times 2} A$$

$$\begin{pmatrix} A'_x \\ A'_y \end{pmatrix}$$

$$\begin{pmatrix} A_x \\ A_y \end{pmatrix}$$

$$\begin{pmatrix} \cos^2 \alpha + e^{i\Delta} \sin^2 \alpha & (1 - e^{i\Delta}) \sin \alpha \cos \alpha \\ (1 - e^{i\Delta}) \sin \alpha \cos \alpha & \sin^2 \alpha + e^{i\Delta} \cos^2 \alpha \end{pmatrix}$$

If the birefringent sheet produces

$\Delta = \frac{\pi}{2}$, it is called a quarter wave plate.

$\Delta = \pi$, half wave plate.

$\Delta = 2\pi$, full wave plate.

M for a quartz wave plate:

$$Q = M = \frac{1}{2} \begin{pmatrix} (1+i) & \pm(1-i) \\ \pm(1-i) & (1+i) \end{pmatrix}$$

$$\Delta = \pi/2$$

$$\alpha = \pm\pi/4$$

$$1+i = 2 e^{-i\pi/4}$$

$$1-i = 2 e^{i\pi/4}$$

$$= \frac{1}{2} \begin{pmatrix} 2 e^{-i\pi/4} & \pm 2 e^{i\pi/4} \\ \pm 2 e^{i\pi/4} & 2 e^{-i\pi/4} \end{pmatrix}$$

$$= e^{i\pi/4} \begin{pmatrix} 1 & \pm e^{i\pi/2} \\ \pm e^{-i\pi/2} & 1 \end{pmatrix}$$

neglect.

$\begin{pmatrix} 1 & \Gamma_i \\ \Gamma_i & 1 \end{pmatrix}$ is the quarter-wave plate transfer function.
