

Lecture 10.

Displacement gradient,  $\nabla \underline{u} = u_{i,j}$

Strain,  $\underline{\varepsilon} = \varepsilon_{ij}$

Rotation,  $\underline{\omega} = \omega_{ij}$

Strain is symmetric  $\varepsilon_{ij} = \varepsilon_{ji}$

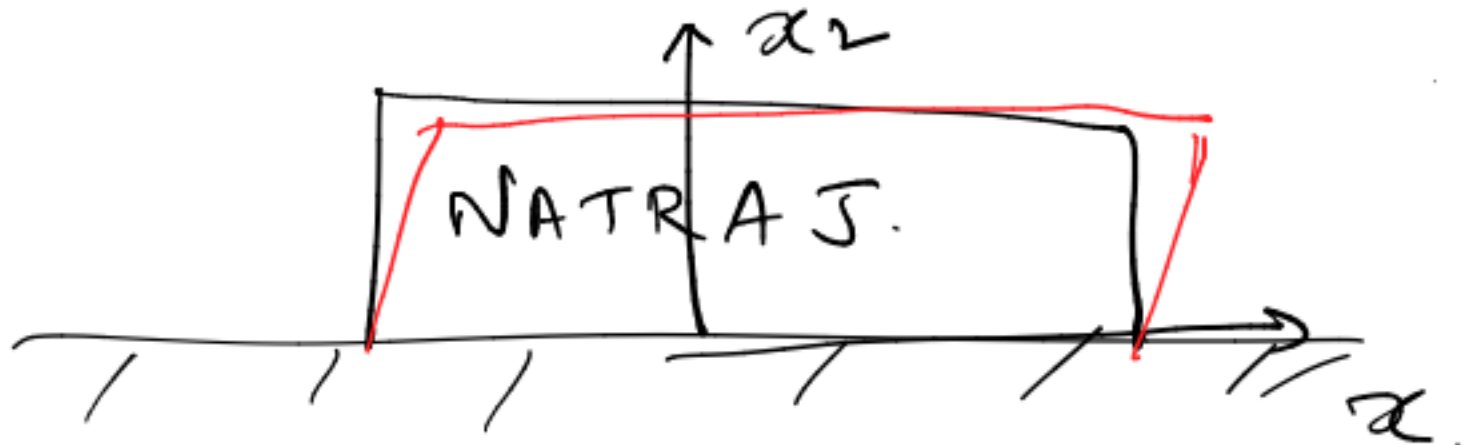
Rotation is skew,  $\omega_{ij} = -\omega_{ji}$

$$\nabla \underline{u} = \underline{\varepsilon} + \underline{\omega}$$

Eg: Simple shear.



Infinite body under simple shear.



$$\tilde{u} = x_2 \tan \theta \quad \tilde{e}_j = c x_2 e_1$$

$$\left( \nabla_{\tilde{u}} \right)_{ij} = \begin{pmatrix} 0 & c & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$\Sigma_{ij} = \begin{pmatrix} 0 & c/2 & 0 \\ c/2 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$\omega_{ij} = \begin{pmatrix} 0 & c/2 & 0 \\ -c/2 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

Q: Suppose there is a displ. field  
for which

$$(\nabla u)_{ij} = \begin{pmatrix} 0 & c/2 & 0 \\ -c/2 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

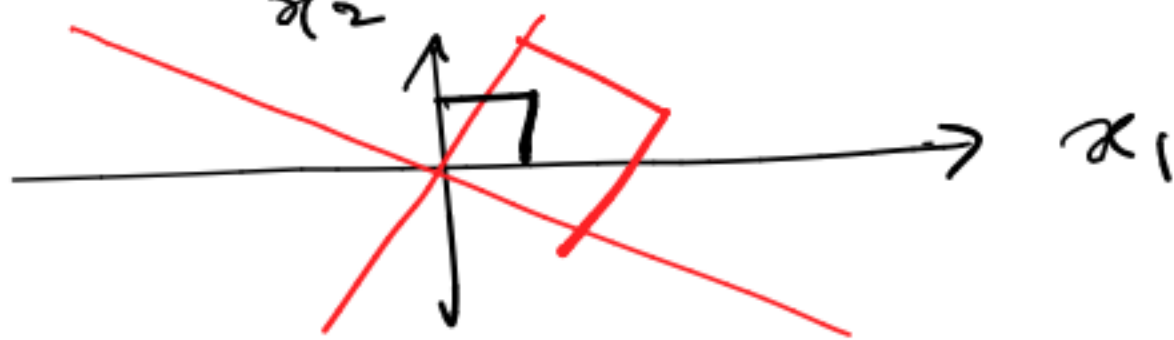
What is the corresponding  $\underline{u}(\underline{x})$ ?

A:  $\frac{\partial u_1}{\partial x_2} = \frac{c}{2} \Rightarrow u_1(x_1, x_2) = \frac{c}{2} x_2 + f(x_1)$

$$\frac{\partial u_2}{\partial x_1} = -\frac{c}{2} \Rightarrow u_2(x_1, x_2) = -\frac{c}{2} x_1 + g(x_2)$$

other  $\frac{\partial u_i}{\partial x_j} = 0$

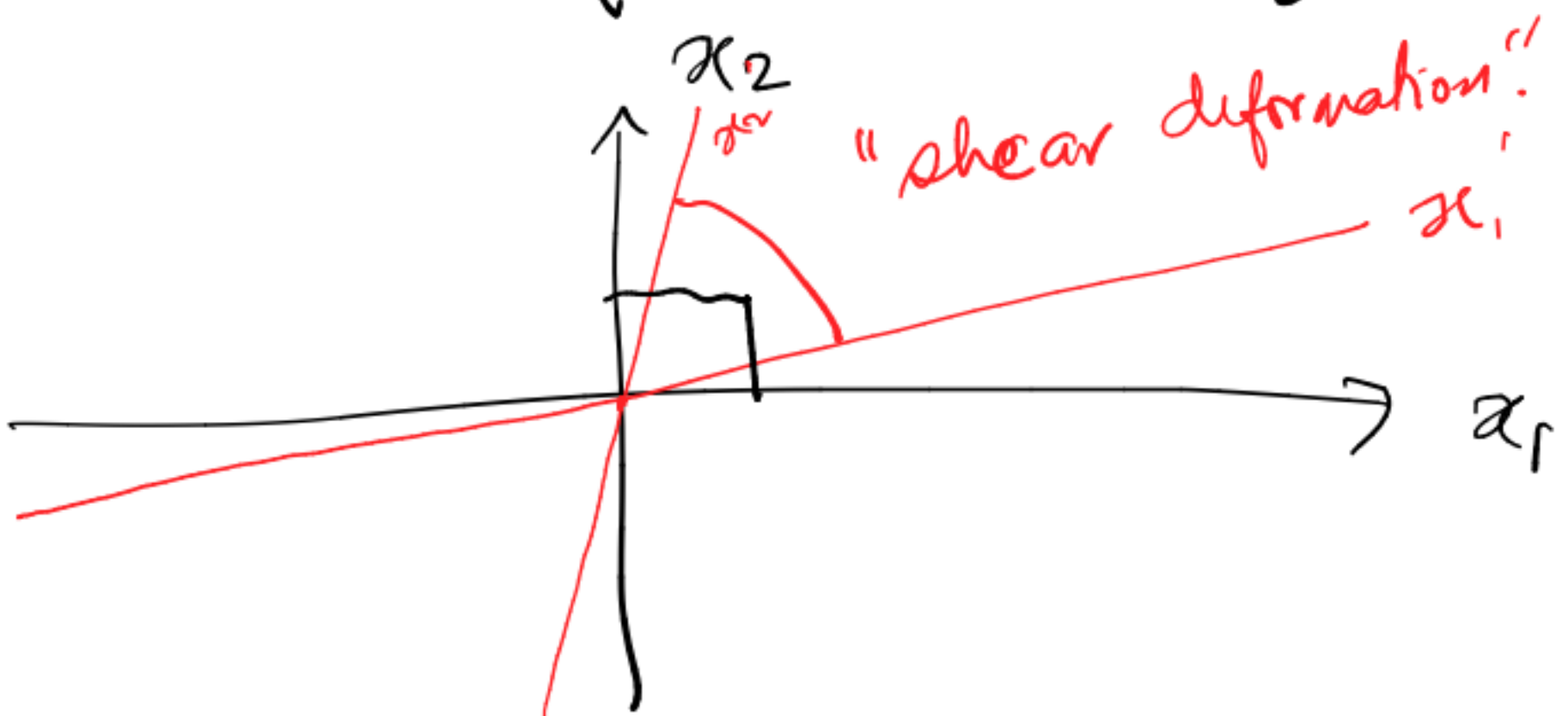
$$\underline{u}(\underline{x}) = \frac{c}{2} x_2 \underline{e}_1 - \frac{c}{2} x_1 \underline{e}_2$$



Because  $\nabla \underline{u} = \underline{\underline{\omega}}$  produces the axes & other lines,  $\underline{\underline{\omega}}$  is called a rotation.

Next

$$(\nabla \underline{u})_{ij} = \begin{pmatrix} 0 & c/2 & 0 \\ c/2 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$



$$\underline{\underline{u}}(\underline{\underline{x}}) = \frac{c}{2} x_2 \underline{\underline{e}}_1 + \frac{c}{2} x_1 \underline{\underline{e}}_2$$

Because angle between the axes changes not  $\nabla \underline{u}$  application,

we have a deformation.

---

§ 2.2, 2.3 Strain transform  
↑  
2-tensor

$$\epsilon'_{ij} = Q_{ip} Q_{jq} \epsilon_{pq}$$

→ apply this formula to 2D

to obtain Mohr's circle for

strain transformation ←  $\epsilon/\omega$

---

2.4 principal strains.

3 invariants

$$I_1 = \text{tr } \epsilon_{ij} = \text{"dilatation"}$$

$I_2$

$I_3 =$

§ 2.5: Decomposition of strain

$$\epsilon_{ij} = \underset{\substack{\uparrow \\ \text{spherical} \\ \text{part} \\ \text{or} \\ \text{hydrostatic} \\ \text{part}}}{\tilde{\epsilon}_{ij}} + \underset{\substack{\uparrow \\ \text{deviatoric} \\ \text{part.}}}{\hat{\epsilon}_{ij}}$$

$$\tilde{\epsilon}_{ij} = \underbrace{\frac{1}{3} \epsilon_{kk}}_{\text{dilatation}} \delta_{ij} \quad \left. \vphantom{\frac{1}{3} \epsilon_{kk}} \right) \text{Captures volume changes.}$$

$$\hat{\epsilon}_{ij} = \epsilon_{ij} - \frac{\epsilon_{kk}}{3} \delta_{ij} \quad \left. \vphantom{\frac{\epsilon_{kk}}{3} \delta_{ij}} \right) \text{distortion of the element}$$



$$\text{of } \left( \underline{\underline{\Sigma}} \right)_{ij} = \begin{pmatrix} 0 & c_2 & 0 \\ c_2 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \quad \Sigma_{ij}' = \underline{\underline{\Sigma}}_{ij}'$$

$$\underline{\underline{\Sigma}}_{ij}' = 0.$$

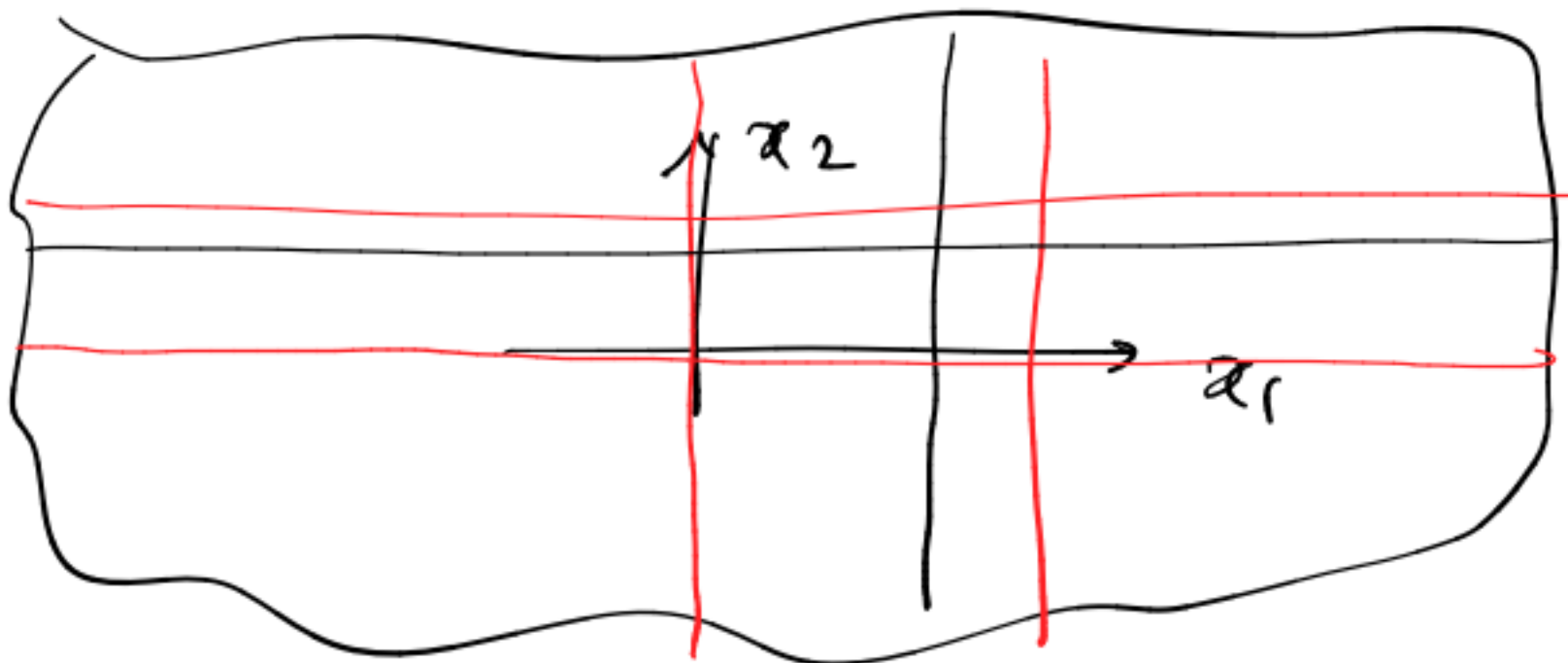
§ 2.7: St. Venant Equations.

Eg.:

$$\underline{\underline{u}}^{(x)} = c x_1 \underline{\underline{e}}_1 + c x_2 \underline{\underline{e}}_2$$

$$\left( \nabla \underline{\underline{u}} \right)_{ij} = \begin{pmatrix} c & 0 & 0 \\ 0 & c & 0 \\ 0 & 0 & 0 \end{pmatrix} = \Sigma_{ij}$$

$$\omega_{ij} \equiv 0.$$



Often we must start w/ the strain field  $\epsilon_{ij}$  & try to calculate  $u_i$  from it.

Q: Can we find the  $u_i$  corresponding to arbitrarily prescribed  $\epsilon_{ij}$ ?

A: No.



